**Communications in Mathematics and Applications** 

Vol. 7, No. 2, pp. 105–113, 2016 ISSN 0975-8607 (online); 0976-5905 (print) Published by RGN Publications



# A Novel Approach for the Stability Analysis of State Dependent Differential Equation Research Article

Sertaç Erman\* and Ali Demir

Department of Mathematics, Kocaeli University, Umuttepe, 41380 Izmit Kocaeli, Turkey **\*Corresponding author:** 106133002@kocaeli.edu.tr

**Abstract.** In this paper, we investigate the stability of a differential equation with state-dependent delay under some conditions on delay term. New necessary and sufficient criterions are elaborated for the asymptotic stability of the differential equations with state dependent delay. Moreover, the asymptotic stability of it is illustrated for a special delay function.

Keywords. Asymptotic stability; State depended delay; Delay differential equation

**MSC.** 37L15

Received: March 14, 2016

Accepted: September 2, 2016

Copyright © 2016 Sertaç Erman and Ali Demir. This is an open access article distributed under the Creative Commons Attribution License, which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited.

### 1. Introduction

State-Dependent delay Differential Equations (SDDE) have a number of applications in various research arias ranging from population dynamics to control theory, see Section 2 in [14] as a review. SDDE is used to make more realistic modelling in the systems whose delay varies according to the internal effects of the system. Therefore they are generating increasing interest from engineers and scientist in recent years. It is shown that the length of time to maturity of Antarctic whales and seals alter according to the state of the population in [13] and it is analyzed by using a mathematical model with SDDE in [1].

SDDE have been investigated for the last five decades. The fundamental theory for local existence and uniqueness theorem for SDDE having Lipschitz continuous initial functions was developed by Drive [8,9] and Driver and Norris [10]. Winston [37] showed that SDDE has a unique solution under some conditions in addition to continuous initial function. There exist some of the earliest studies on SDDE in [2,6,32]. Moreover, lots of theoretical and numerical

analysis of SDDE have been done so far, [1, 3–5, 11, 14–16, 18, 20, 22–30, 33–35]. Especially [36] can be seen as a review in order to have much more detail about DDE and SDDE and researches on them.

In this study, we consider the following type of SDDE

$$u'(t) = -A_0 u(t) - A_1 u(t - \tau(t, u(t))), \tag{1}$$

where  $A_0$ ,  $A_1 \in \mathbb{R}$  and  $\tau(t, u(t)) > 0$ ,  $\tau(t, 0) \neq 0$  for all  $t \in \mathbb{R}^+$ . Since the variation of solution can play an important role in the variation of delay, hence its stability, we use the following characteristic equation

$$g(\lambda) = \lambda + A_0 + A_1 e^{-\lambda \tau(t, e^{\lambda t})} = 0$$
<sup>(2)</sup>

in order to analyze stability of solution of equation (1).

In the general case, the characteristic roots  $\lambda_j$ , j = 1, 2, ..., of equation (1) are obtained by solving the characteristic equation (2) where  $\lambda_j$  is a complex number. If the characteristic roots have negative real parts, i.e.,  $\operatorname{Re}(\lambda_j) < 0$  for all j = 1, 2, ... then the solution of (1) is asymptotically stable and if at least one of the characteristic roots have positive real parts, i.e.,  $\operatorname{Re}(\lambda_j) > 0$  for some j = 1, 2, ... then the solution of (1) is unstable.

It is attempted to determine the stability and instability regions of the system in parameter space  $(A_0, A_1)$  by using D-partition method. The method is originated from paper [31]. Moreover, the analysis by this method are conducted in [12,17,19,21] and [7]. We consider the characteristic equation  $g(\lambda, A_0, A_1)$  in two parameters for equation (1). D-partition method is based on fact that the roots of the characteristic equation are continuous functions of the parameters  $A_0$  and  $A_1$ . Varying the parameters,  $\lambda_j$  change continuously in complex plane and at least one  $\lambda_j$  crosses the imaginary axis at the point where the stability changes. In this method parameter space is divided into subregions by the hypersurfaces called the D-curves. The points of the D-curves correspond to pure imaginary roots or zero root of the characteristic equation. Moreover, the characteristic equation has the same number of roots with positive real part in each subregion in the parameter space determined by the D-curves. Thus, finding specific point at which the solution of equation (1) is stable, it is enough to find the stability region, including this point.

In order to obtain D-curves, pure imaginary number  $\lambda = i\omega$  is substituted in characteristic equation  $g(\lambda, A_0, A_1)$ . Equating to zero the real and imaginary parts, we have

$$U(\omega, A_0, A_1) = \text{Re}(g(i\omega, A_0, A_1)) = 0,$$
(3)

$$V(\omega, A_0, A_1) = \text{Im}(g(i\omega, A_0, A_1)) = 0.$$
(4)

Hence, by making use of (3) and (4), parametric equations can be rewritten in the following form

$$A_0 = A_0(\omega), \quad A_1 = A_1(\omega)$$

where  $\omega$  is a parameter, ranging from  $-\infty$  to  $\infty$ . These curves and singular solutions of equations (3) and (4) constitute D-curves.

In Section 2, we establish necessary and sufficient conditions for the stability of the solution of equation (1) by using D-partition method. We illustrate these results in Section 3.

## 2. Necessary and Sufficient Conditions for the Stability of SDDE

It is clear that if  $\tau(t, u(t))$  is a linear function with respect to u(t), then  $\operatorname{Re}(\tau(t, e^{i\omega t}))$  is an even function and  $\operatorname{Im}(\tau(t, e^{i\omega t}))$  is an odd function with respect to  $\omega$ . In this section, we consider equation (1) for  $\tau(t, e^{i\omega t}) = \tau_1(t, \omega) + i\tau_2(t, \omega)$  such that  $\tau_1(t, 0) \neq 0$ ,  $\tau_1(t, \omega)$  is an even function and  $\tau_2(t, \omega)$  is an odd function with respect to  $\omega$ .

The characteristic equation

$$g(\lambda) = \lambda + A_0 + A_1 e^{-\lambda \tau(t, e^{\lambda t})} = 0.$$
(5)

As a part of the D-partition method, we have

$$C_*: A_0 + A_1 = 0, \quad \text{for } \lambda = 0$$
 (6)

this straight line is a line forming the boundary of the D-partition and is denoted by  $C_*$ . Substituting  $\lambda = i\omega$  and equating to zero the real and imaginary parts in characteristic equation (5), we find following equations

$$A_0 + A_1 e^{\omega \tau_2(t,\omega)} \cos(\omega \tau_1(t,\omega)) = 0,$$
(7)

$$\omega + A_1 e^{\omega \tau_2(t,\omega)} \sin(\omega \tau_1(t,\omega)) = 0.$$
(8)

Solving the above equations for  $A_0$  and  $A_1$ , following parametric curve equations are obtained

$$A_0(\omega) = -\frac{\omega \cos(\omega \tau_1(t,\omega))}{\sin(\omega \tau_1(t,\omega))},\tag{9}$$

$$A_1(\omega) = \frac{\omega}{e^{\omega \tau_2(t,\omega)} \sin(\omega \tau_1(t,\omega))}.$$
(10)

Since  $A_0(\omega)$  and  $A_1(\omega)$  are an even with respect to  $\omega$ , it is sufficient to take  $\omega \in (0,\infty)$ . When  $\omega$  ranges from 0 to  $\infty$ , equations (9)-(10) define the D-curves for each  $t \in \mathbb{R}^+$ . The equations (9)-(10) have singularity for  $\omega \tau_1(t,\omega) = k\pi$ ,  $k = 0, 1, 2, \ldots$ . Thus, we introduce intervals  $J_k = (k\pi, (k+1)\pi)$  and denote by  $C_k$  the curve in the parameter space  $(A_0, A_1)$  for  $\omega \tau_1(t, \omega) \in J_k$ .

To analyze the curves  $C_k$  under assumptions  $\tau_1(t, \omega) \le h_1$  and  $\omega \tau_2(t, \omega) \le h_2$  where  $h_1$  and  $h_2$  are non-negative real numbers, we firstly consider the following equation

$$\widetilde{A}_1(\omega, T_2) = \frac{\omega}{T_2 \sin(\omega \tau_1(t, \omega))}$$
(11)

where  $T_2 \in (0, e^{h_2}]$  is a real number. Equations (9)-(11) define a family of curves since  $T_2$  is not a constant. Holding  $T_2$  fixed, these define  $\tilde{A}_1(\omega, T_2)$  as function of  $\omega$ , providing a parametric representation of a curve. Different values of  $T_2$  give different curves in the family. We denote the family of curves by  $\tilde{C}_k(T_2)$  for  $\omega \tau_1(t, \omega) \in J_k$ .

**Proposition 1.** If  $\omega \tau_2(t, \omega) \leq h_2$ , the curve  $\tilde{C}_{2k}(e^{h_2})$  lies below the curve  $C_{2k}$  in parameter space  $(A_0, A_1)$  for k = 0, 1, 2, ...

*Proof.* For every  $\omega \tau_1(t, \omega) \in J_{2k}$ , equations (9)-(10) give a point  $L(A_0, A_1)$  on the curve  $C_{2k}$  and equations (9)-(11) give a point  $\tilde{L}(A_0, \tilde{A}_1)$  on the curve  $\tilde{C}_{2k}(e^{h_2})$ . Since  $e^{\omega \tau_2(t,\omega)} \leq e^{h_2}$  when  $\omega \tau_2(t,\omega) \leq h_2$  and  $A_1(\omega) > 0$ ,  $\tilde{A}_1(\omega, T_2) > 0$  for each  $\omega \tau_1(t,\omega) \in J_{2k}$ , we have  $A_1 > \tilde{A}_1$ .

**Proposition 2.** If  $\omega \tau_2(t, \omega) \leq h_2$ , the curve  $\widetilde{C}_{2k+1}(e^{h_2})$  lies above the curve  $C_{2k+1}$  in parameter space  $(A_0, A_1)$  for k = 0, 1, 2, ...

*Proof.* It is similar to the proof of Proposition 1.

**Lemma 1.** The curves  $\tilde{C}_k(T_2)$  do not intersect each other for k = 0, 1, 2, ....

*Proof.* Suppose that there exist an intersection point. It means that, there exist  $\omega_1 \neq \omega_2 \in \mathbb{R}^+$  such that  $A_0(\omega_1, T_2) = A_0(\omega_2, T_2)$  and  $\widetilde{A}_1(\omega_1, T_2) = \widetilde{A}_1(\omega_2, T_2)$ . These equalities imply that

$$\frac{\omega_1}{T_2 \sin(\omega_1 \tau_1(t,\omega_1))} = \frac{\omega_2}{T_2 \sin(\omega_2 \tau_1(t,\omega_2))}, \quad \frac{\omega_1 \cos(\omega_1 \tau_1(t,\omega_1))}{\sin(\omega_1 \tau_1(t,\omega_1))} = \frac{\omega_2 \cos(\omega_2 \tau_1(t,\omega_2))}{\sin(\omega_2 \tau_1(t,\omega_2))}$$
(12)

from equation (9) and (11). For  $n \in \mathbb{N}$ ,  $\omega_1 \tau_1(t, \omega_1) \neq \omega_2 \tau_1(t, \omega_2) + 2n\pi$  is obtained from the left equality in (12) because of  $\omega_1 \neq \omega_2$ . In addition, left and right equalities in (12) lead to  $\cos(\omega_1 \tau_1(t, \omega_1)) = \cos(\omega_2 \tau_1(t, \omega_2))$  which is a contradiction.

**Lemma 2.** The following limits are satisfied for k = 1, 2, ...

$$\lim_{\substack{\omega\tau_1(t,\omega)\to(2k\pi)^+ \\ \omega\tau_1(t,\omega)\to(2k-1)\pi)^-}} A_0(\omega) = \lim_{\substack{\omega\tau_1(t,\omega)\to(2k\pi)^+ \\ \omega\tau_1(t,\omega)\to((2k-1)\pi)^-}} A_1(\omega,h) = +\infty,$$
(13)

and

$$\lim_{\substack{\omega\tau_1(t,\omega)\to(2k\pi)^-\\ \lim_{\omega\tau_1(t,\omega)\to((2k-1)\pi)^+}}} A_0(\omega) = \lim_{\substack{\omega\tau_1(t,\omega)\to(2k\pi)^-\\ \lim_{\omega\tau_1(t,\omega)\to((2k-1)\pi)^+}}} A_1(\omega,h) = -\infty,$$
(14)

Then, we consider the following family of curves denoted by  $\overline{C}_k(T_1)$ 

$$\overline{C}_{k}(T_{1}): \begin{cases} \overline{A}_{0}(\omega, T_{1}) = -\frac{\omega \cos(\omega T_{1})}{\sin(\omega T_{1})}, & T_{1} \in (-\infty, h_{1}], \\ \\ \overline{A}_{1}(\omega, T_{1}) = \frac{\omega}{e^{h_{2}} \sin(\omega T_{1})}, & \omega T_{1} \in J_{k}. \end{cases}$$
(15)

**Proposition 3.** If  $\tau_1(t,\omega) \leq h_1$  and  $\omega \tau_2(t,\omega) \leq h_2$ , the curve  $\overline{C}_{2k}(h_1)$  lies below the curve  $C_{2k}$  in parameter space  $(A_0, A_1)$  for  $\omega h_1 \in J_{2k}$ , k = 0, 1, 2, ...

*Proof.*  $A_0(\omega) < \overline{A}_0(\omega, h_1)$  for all  $\omega h_1 \in J_k$ , since the following partial derivative of  $\overline{A}_0(\omega, T_1)$  with respect to  $T_1$ 

$$\frac{\partial \overline{A}_0}{\partial T_1} = \frac{\omega^2}{\sin^2(\omega T_1)} > 0, \quad \forall \ \omega T_1 \in J_k.$$

Moreover, taking the derivative of  $\overline{A}_1(\omega, T_1)$  with respect to  $T_1$ , we obtain

$$\frac{\partial A_1}{\partial h} = \frac{-\omega^2 \cos(\omega T_1)}{e^{h_2} \sin^2(\omega T_1)}$$

and  $\overline{A}_1(\omega, T_1)$  is a monotone decreasing function for  $\omega T_1 \in \left(2k\pi, \frac{(2k+1)\pi}{2}\right)$ . Therefore, the point  $\overline{L}(\overline{A}_0(\omega, h_1), \overline{A}_1(\omega, h_1))$  lies below the point  $\widetilde{L}(A_0(\omega), \widetilde{A}_1(\omega, e^{h_2}))$  for  $\omega h_1 \in \left(2k\pi, \frac{(2k+1)\pi}{2}\right)$ . Because of Lemma 1 and limits (13),  $\overline{C}_{2k}(h_1)$  lies below the curve  $\widetilde{C}_{2k}(e^{h_2})$  for  $\omega h_1 \in J_{2k}$ . Hence,  $\overline{C}_{2k}(h_1)$  lies below the curve  $C_{2k}$  for  $\omega h_1 \in J_{2k}$ .

**Proposition 4.** If  $\tau_1(t,\omega) \leq h_1$  and  $\omega \tau_2(t,\omega) \leq h_2$ , the curve  $\overline{C}_{2k+1}(h_1)$  lies above the curve  $C_{2k+1}$  in parameter space  $(A_0, A_1)$  for  $\omega h_1 \in J_{2k+1}$ , k = 0, 1, 2, ...

*Proof.* It is similar to the proof of Proposition 3.

**Lemma 3.** If  $e^{h_2} \leq 1$ , the curve  $\overline{C}_0(h_1)$  intersects  $C_*$  exactly once at  $\left(-\frac{1}{h_1}, \frac{1}{e^{h_2}h_1}\right)$  and  $\overline{C}_k(h_1)$  do not intersect  $C_*$  for k = 1, 2, ...

*Proof.* Intersection of  $\overline{C}_0(h_1)$  and  $C_*$  is obvious from the following limit point

$$\left(\lim_{\omega\to 0} \overline{A}_0(\omega,h_1), \lim_{\omega\to 0} \overline{A}_1(\omega,h_1)\right) = \left(-\frac{1}{h_1}, \frac{1}{e^{h_2}h_1}\right).$$

Assume that  $\overline{C}_k(h_1)$  and  $C_*$  has intersection points, then there exist  $\omega h_1 \in J_k$  for equations (15) which satisfies equation (6). By using equations (15) in equation (6), we have

$$\frac{\omega\cos(\omega h_1)}{\sin(\omega h_1)} = \frac{\omega}{e^{h_2}\sin(\omega h_1)}$$

which has no solution  $\omega h_1 \in J_k$  for k = 0, 1, 2, ... and this contradicts with our assumption.

**Lemma 4.** If  $e^{h_2} > 1$ , the curves  $\overline{C}_k(h_1)$  intersect  $C_*$  at point  $(\overline{A}_0(\omega_0, h_1), \overline{A}_1(\omega_0, h_1))$  where  $\omega_0$  is the root of  $\omega_0 = \frac{1}{h_1} \arccos\left(\frac{1}{e^{h_2}}\right)$  such that  $\omega_0 h_1 \in J_k$ .

*Proof.* A straightforward computation shows that the corresponding point  $(\overline{A}_0(\omega_0, h_1), \overline{A}_1(\omega_0, h_1))$  lies on the line  $C_*$ .

**Lemma 5.** The curve  $\overline{C}_k(h_1)$  intersects the line  $A_0 = 0$  exactly once. Moreover, the intersection point  $(0, P_k)$  satisfies the following inequalities

$$\begin{array}{ll} P_k < P_{k+2}, & \mbox{for } k = 2n, \ n \in \mathbb{N}, \\ P_{k+2} < P_k, & \mbox{for } k = 2n+1, \ n \in \mathbb{N}. \end{array}$$

*Proof.* When  $\omega h_1 \in J_k$ , the equation  $\overline{A}_0(\omega, h_1) = 0$  implies  $\omega = \frac{\pi + 2k\pi}{2h_1}$ . Hence,

$$P_{k} = \begin{cases} \frac{\pi + 2k\pi}{2e^{h_{2}}h_{1}} & \text{ for } k = 2n, \ n \in \mathbb{N} \\ \\ -\frac{\pi + 2k\pi}{2e^{h_{2}}h_{1}} & \text{ for } k = 2n + 1, \ n \in \mathbb{N} \end{cases}$$

is obtained by substituting  $\omega = \frac{\pi + 2k\pi}{2h_1}$  in  $\overline{A}_1(\omega, h_1)$ . This completes the proof.

**Lemma 6.** The following limits are satisfied for k = 1, 2, ...

$$\lim_{\omega h_1 \to \left(\frac{(2k-1)\pi}{h}\right)^-} \overline{A}_0(\omega, h_1) = \lim_{\omega h_1 \to \left(\frac{(2k-1)\pi}{h}\right)^-} \overline{A}_1(\omega, h_1) = \lim_{\omega h_1 \to \left(\frac{2k\pi}{h}\right)^-} \overline{A}_0(\omega, h_1)$$
$$= \lim_{\omega h_1 \to \left(\frac{2k\pi}{h}\right)^+} \overline{A}_1(\omega, h_1) = +\infty$$
$$\lim_{\omega h_1 \to \left(\frac{(2k-1)\pi}{h}\right)^+} \overline{A}_0(\omega, h_1) = \lim_{\omega h_1 \to \left(\frac{(2k-1)\pi}{h}\right)^+} \overline{A}_1(\omega, h_1) = \lim_{\omega h_1 \to \left(\frac{(2k\pi)\pi}{h}\right)^+} \overline{A}_0(\omega, h_1)$$

$$\lim_{\omega h_1 \to \left(\frac{(2k-1)\pi}{h}\right)^+} A_0(\omega, h_1) = \lim_{\omega h_1 \to \left(\frac{(2k-1)\pi}{h}\right)^+} A_1(\omega, h_1) = \lim_{\omega h_1 \to \left(\frac{2k\pi}{h}\right)^-} \overline{A}_1(\omega, h_1) = -\infty$$

Communications in Mathematics and Applications, Vol. 7, No. 2, pp. 105–113, 2016

**Theorem 1.** Suppose that  $\tau_1(t,0) \neq 0$ ,  $\tau_1(t,\omega)$  is an even function,  $\tau_2(t,\omega)$  is an odd function with respect to  $\omega$ . Moreover,  $\tau_1(t,\omega) \leq h_1$ ,  $\omega \tau_2(t,\omega) \leq h_2$  where  $h_1$  and  $h_2$  are non-negative real numbers and  $e^{h_2} \leq 1$ . The solution of equation (1) is asymptotically stable if the following conditions are satisfied:

(i) 
$$-\frac{1}{h_1} < A_0$$
  
(ii)  $-A_0 < A_1 < \frac{\omega}{e^{h_2} \sin(\omega h_1)}$  where  $\omega$  is the root of  $A_0 = -\frac{\omega \cos(\omega h_1)}{\sin(\omega h_1)}$  such that  $\omega h_1 \in J_0$ .

*Proof.* When  $A_0 > 0$  and  $A_1 = 0$ , the solution of equation (1) is clearly asymptotically stable. The stability region which includes half line  $A_0 > 0$  and  $A_1 = 0$ , lies above  $C_*$  and below  $\overline{C}_0(h_1)$  as a result of Proposition 3, Proposition 4, Lemma 3, Lemma 5 and Lemma 6. The conditions (i)-(ii) are algebraic representation of this region in parameter space  $(A_0, A_1)$ .

**Theorem 2.** Suppose that  $\tau_1(t,0) \neq 0$ ,  $\tau_1(t,\omega)$  is an even function and  $\tau_2(t,\omega)$  is an odd function with respect to  $\omega$ . Moreover,  $\tau_1(t,\omega) \leq h_1$ ,  $\omega \tau_2(t,\omega) \leq h_2$  where  $h_1$  and  $h_2$  are non-negative real numbers and  $e^{h_2} > 1$ . The solution of equation (1) is asymptotically stable if the following conditions are satisfied:

(iii)  $-\frac{1}{h_1} < A_0 \text{ or } \overline{A}_0(\omega, h_1) < A_0 \text{ where } \omega \text{ is the root of } \omega = \frac{1}{h_1} \arccos\left(\frac{1}{e^{h_2}}\right) \text{ such that } \omega h_1 \in J_0$ (iv)  $-A_0 < A_1 < \frac{\omega}{e^{h_2} \sin(\omega h_1)} \text{ where } \omega \text{ is the root of } A_0 = -\frac{\omega \cos(\omega h_1)}{\sin(\omega h_1)} \text{ such that } \omega h_1 \in J_0$ (v)  $\frac{\omega}{e^{h_2} \sin(\omega h_1)} < A_1 \text{ where } \omega \text{ is the root of } A_0 = -\frac{\omega \cos(\omega h_1)}{\sin(\omega h_1)} \text{ such that } \omega h_1 \in J_1.$ 

*Proof.* The stability region which includes half line  $A_0 > 0$  and  $A_1 = 0$ , lies among  $C_*$ ,  $\overline{C}_1(h_1)$  and  $\overline{C}_0(h_1)$  because of Proposition 3, Proposition 4, Lemma 4, Lemma 5 and Lemma 6. The conditions (iii), (iv) and (v) are algebraic representation of this region in parameter space  $(A_0, A_1)$ .

**Theorem 3.** Suppose that  $\tau_1(t,0) \neq 0$ ,  $\tau_1(t,\omega)$  is an even function and  $\tau_2(t,\omega)$  is an odd function with respect to  $\omega$  and  $\omega \tau_2(t,\omega) \leq h_2$  where  $h_2$  is a non-negative real number. The solution of equation (1) is asymptotically stable, if the following condition is satisfied:

(vi)  $A_0 \leq |A_1 e^{h_2}|$ .

*Proof.* It is obvious from (7) that,  $A_0 \leq |A_1e^{h_2}|$  for all  $\omega \in J_k$ . Therefore there is no D-curve in the region described by (vi). Moreover, the half line  $A_0 > 0$  and  $A_1 = 0$  on which the solution equation (1) is asymptotically stable, is in this region.

**Theorem 4.** Suppose that  $A_1 \neq 0$ ,  $\tau_1(t,0) \neq 0$ ,  $\tau_1(t,\omega)$  is an even function and  $\tau_2(t,\omega)$  is an odd function with respect to  $\omega$ . If  $\omega \tau_2(t,\omega)$  does not have an upper bound, then the solution of equation (1) is not stable.

*Proof.* It follows from (10) that we have

 $\lim_{\omega\tau_2(t,\omega)\to\infty} A_1(\omega) = \lim_{\omega\tau_2(t,\omega)\to\infty} \frac{\omega}{e^{\omega\tau_2(t,\omega)}\sin(\omega\tau_1(t,\omega))} = 0.$ 

Thus, D-curves tend to half line  $A_0 > 0$  and  $A_1 = 0$ , when  $\omega$  ranges from 0 to  $\infty$ .

# 3. The Stability of Equation (1) with Delay Term $\tau(t, u(t)) = \frac{au(t)+b}{cu(t)+d}$

We consider the stability of equation (1) with delay term  $\tau(t, u(t)) = \frac{au(t)+b}{cu(t)+d}$  where *a*, *b*, *c* and *d* are positive real numbers.

The Möbius transformation can be rewritten as follows

$$\tau(z) = \frac{az+b}{cz+d} = \frac{bc-ad}{c} \frac{1}{cz+d} + \frac{a}{c}$$

where  $z \in \mathbb{C}$ . Hence,

$$\tau_1(t,\omega) = \frac{a}{c}, \quad \tau_2(t,\omega) = 0$$

when bc - ad = 0. It follows from Theorem 1 that if bc - ad = 0 and the following conditions are satisfied:

(i) 
$$-\frac{c}{a} < A_0$$

(ii)  $-A_0 < A_1 < \frac{\omega}{\sin\left(\omega\frac{a}{c}\right)}$  where  $\omega$  is the root of  $A_0 = -\frac{\omega\cos\left(\frac{\omega a}{c}\right)}{\sin\left(\frac{\omega a}{c}\right)}$  such that  $\frac{\omega a}{c} \in J_0$ 

then the solution of equation (1) with delay term  $\tau(t, u(t)) = \frac{au(t)+b}{cu(t)+d}$  is asymptotically stable.

#### 4. Conclusion

In this study, stability conditions are given in terms of the coefficients in equation (1) under some conditions on the delay function  $\tau(t, u(t))$  in Theorem 1, Theorem 2 and Theorem 3. Moreover, it is proved that the condition  $\tau_2(t, \omega) \le h_2$  given in Theorem 4, is the necessary condition for the stability of the solution.

In literature, state dependent delays are linearized heuristically by freezing at a constant solution in order to investigate the stability of SDDE. Heuristic linearization is applied taking  $\tau(t, u(t)) = \tau(t, 0)$  in equation (1) in case of zero solution. It is show that, the stability analysis under the condition bc - ad = 0 is the same as the one which is obtained by linearization at zero in Section 3.

#### **Competing Interests**

The authors declare that they have no competing interests.

#### **Authors' Contributions**

All the authors contributed significantly in writing this article. The authors read and approved the final manuscript.

#### References

[1] W.G. Aiello, H.I. Freedman and J. Wu, Analysis of a model representing stage-structured population growth with statedependent time delay, *SIAM J. Appl. Math.* **52** (1992), 855–869.

- [2] W. Alt, Some periodicity criteria for functional differential equations, *Manuscripta Math.* 23 (1978), 295–318.
- [3] O. Arino, K.P. Hadeler and M.L. Hbid, Existence of periodic solutions for delay differential equations with state-dependent delay, J. Differ. Equ. 144 (1998), 263–301.
- [4] I. Chueshov and A. Rezounenko, Dynamics of second order in time evolution equations with state-dependent delay, *Nonlinear Analysis* (2015), 126–149.
- [5] K. Cooke and W. Huang, On the problem of linearization for state-dependent delay differential equations, *Proc. Am. Math. Soc.* **124** (1996), 1417–1426.
- [6] K.L. Cooke, Asymptotic theory for the delay-differential equation u(t) = -au(t r(u(t))), J. Math. Anal. Appl. 19 (1967), 160–173.
- [7] O. Diekmann, S.A. van Gils, S.M.V. Lunel and H.O. Walther, *Delay Equations, Functional, Complex and Nonlinear Analysis*, Springer, New York (1995).
- [8] R.D. Driver, A two-body problem of classical electrodynamics: the one-dimensional case, Ann. Physics 21 (1963), 122–142.
- [9] R.D. Driver, Existence theory for a delay-differential system, *Contrib. Differential Equations* 1 (1963), 317–336.
- [10] R.D. Driver and M.J. Norris, Note on uniqueness for a one-dimensional two-body problem of classical electrodynamics, Ann. Physics 42 (1967), 347–351.
- [11] M. Eichmann, A local Hopf bifurcation theorem for differential equations with state-dependent delays, Doctoral dissertation, Gießen (2006).
- [12] L.E. Elsgolts and S.B. Norkin, Introduction to the Theory and Application of Differential Equations with Deviating Arguments, Academic Press, London (1973).
- [13] R. Gambell, Birds and mammals-Antarctic whales in Antarctica, in W.N. Bonner and D.W.H. Walton (eds.), Pergamon Press, New York, 223–241 (1985).
- [14] F. Hartung, T. Krisztin, H.O. Walther and J. Wu, Functional differential equations with statedependent delays: theory and applications, in A. Canada, P. Drabek and A. Fonda (eds.), *Handbook* of Differential Equations: Ordinary Differential Equations, Vol. III, Elsevier/North-Holland, Amsterdam, 435-545 (2006).
- [15] F. Hartung, Differentiability of solutions with respect to the initial data in differential equations with state-dependent delay, J. Dyn. Differ. Equ. 23 (2011), 843–884.
- [16] Q. Hu and J. Wu, Global Hopf bifurcation for differential equations with state-dependent delay, J. Differ. Equ. 248 (2010), 2081–2840.
- [17] T. Insperger and G. Stépán, Semi-Discretization Stability and Engineering Applications for Time-Delay Systems, Springer, New York (2011).
- [18] B. Kennedy, Multiple periodic solutions of an equation with state-dependent delay, J. Dyn. Differ. Equ. 26 (2011), 1–31.
- [19] V.B. Kolmanovskii and V.R. Nosov, *Stability of Functional Differential Equations*, Academic Press, London (1986).
- [20] G. Kozyreff and T. Erneux, Singular Hopf bifurcation in a differential equation with large statedependent delay, Proc. R. Soc. A 470, doi: 10.1098/rspa.2013.0596 (2014).
- [21] A.M. Krall, *Stability Techniques for Continuous Linear Systems*, Gordon and Breach, New York (1967).

- [22] T. Krisztin and O. Arino, The 2-dimensional attractor of a differential equation with state-dependent delay, *J. Dyn. Differ. Equ.* 13 (2001), 453–522.
- [23] T. Krisztin, An unstable manifold near a hyperbolic equilibrium for a class of differential equations with state-dependent delay, *Discrete Contin. Dyn. Syst.* 9 (2003), 993–1028.
- [24] Y. Kuang and H.L. Smith, Slowly oscillating periodic solutions of autonomous state-dependent delay differential equations, *Nonlinear Anal. Theory Methods Appl.* 19 (1992), 855–872.
- [25] P. Magal and O. Arino, Existence of periodic solutions for a state-dependent delay differential equation, J. Differ. Equ. 165 (2000), 61–95.
- [26] J. Mallet-Paret, R.D. Nussbaum and P. Paraskevopoulos, Periodic solutions for functional differential equations with multiple state-dependent time lags, *Topol. Methods Nonlinear Anal.* 3 (1994), 101–162.
- [27] J. Mallet-Paret and R.D. Nussbaum, Boundary layer phenomena for differential-delay equations with state-dependent time-lags, *I. Arch. Ration Mech. Anal.* **120** (1992), 99–146.
- [28] J. Mallet-Paret and R.D. Nussbaum, Boundary layer phenomena for differential-delay equations with state-dependent time-lags: II, J. Reine Angew. Math. 477 (1996), 129–197.
- [29] J. Mallet-Paret and R.D. Nussbaum, Boundary layer phenomena for differential-delay equations with state-dependent time-lags: III, J. Differ. Equ. 189 (2003), 640–692.
- [30] J. Mallet-Paret and R.D. Nussbaum, Superstability and rigorous asymptotics in singularly perturbed state-dependent delay-differential equations, J. Differ. Equ. 250 (2011), 4037–4084.
- [31] J.I. Neimark, D-subdivision and spaces of quasi-polynomials (in Russian), *Prikl Mat. Mekh.* 13 (1949), 349–380.
- [32] R.D. Nussbaum, Periodic solutions of some nonlinear autonomous functional differential equations, Ann. Mat. Pura Appl. 101 (1974), 263–306.
- [33] H.O. Walther, A periodic solution of a differential equation with state-dependent delay, J. Differ. Equ. 244 (2008), 1910–1945.
- [34] J. Sieber, Finding periodic orbits in state-dependent delay differential equations as roots of algebraic equations, *Discrete Contin. Dyn. Syst. Ser. A* 32 (2012), 2607–2651.
- [35] E. Stumpf, On a differential equation with state-dependent delay: a global center-unstable manifold connecting an equilibrium and a periodic orbit, *J. Dyn. Differ. Equ.* **24** (2012), 197–248.
- [36] H.O. Walther, Topics in delay differential equation, Jahresber Dtsch. Math-Ver. 116 (2014), 87-114.
- [37] E. Winston, Uniqueness of solutions of state dependent delay differential equations, J. Math. Anal. Appl. 47 (1974), 620–625.