# Lower Order Eigenvalues of the Schrödinger Operator* 

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#### Abstract

Making use of the method introduced by Brands in [2], we consider lower order eigenvalues of the Schrödinger operator in Euclidean domains. We extend an estimate on eigenvalues obtained by Ashbaugh and Benguria in [1].


Keywords. Membrane eigenvalue; Schrödinger operator; Rayleigh-Ritz inequality
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## 1. Introduction

Let $\Omega$ be a bounded domain in an $n$-dimensional Euclidean space $\mathbb{R}^{n}$ with smooth boundary $\partial \Omega$. The eigenvalue problem

$$
\begin{cases}-\Delta u=\lambda u, & \text { in } \Omega  \tag{1.1}\\ u=0, & \text { on } \partial \Omega\end{cases}
$$

is called the fixed membrane problem. Let $0<\lambda_{1}<\lambda_{2} \leq \lambda_{3} \leq \cdots \rightarrow+\infty$ denote the successive eigenvalues for (1.1), where each eigenvalue is repeated according to its multiplicity. In the case of $n=2$, Payne-Pólya-Weinberger [5] proved

$$
\begin{equation*}
\frac{\lambda_{2}+\lambda_{3}}{\lambda_{1}} \leq 6 \tag{1.2}
\end{equation*}
$$

Subsequently, in 1964, Brands [2] sharpen (1.2) to

$$
\begin{equation*}
\frac{\lambda_{2}+\lambda_{3}}{\lambda_{1}} \leq 5+\frac{\lambda_{1}}{\lambda_{2}} \tag{1.3}
\end{equation*}
$$

[^0]In 1993, for general dimensions $n \geq 2$, Ashbaugh and Benguria [1] proved (see the inequality (6.10) in [1])

$$
\begin{equation*}
\frac{\lambda_{2}+\lambda_{3}+\cdots+\lambda_{n+1}}{\lambda_{1}} \leq n+3+\frac{\lambda_{1}}{\lambda_{2}} . \tag{1.4}
\end{equation*}
$$

Recently, the inequality (1.4) has been extended to some Riemannian manifolds, see [6, 3, 4] and the references therein.

In this note, we consider eigenvalue problem of the following Schrödinger operator

$$
\begin{cases}(-\Delta+V) u=\lambda u, & \text { in } \Omega  \tag{1.5}\\ u=0, & \text { on } \partial \Omega\end{cases}
$$

where $V$ is a continuous bounded function on $\bar{\Omega}$. Using the method of Brands [2], we study the eigenvalue problem (1.5) for general dimensions $n \geq 2$ and extend the inequality (1.4) as follows:

Theorem. Let $\lambda_{i}$ be the $i$-th eigenvalue of the eigenvalue problem (1.5). Then

$$
\begin{equation*}
\frac{\lambda_{2}+\lambda_{3}+\cdots+\lambda_{n+1}}{\lambda_{1}} \leq n+\frac{(M+1)(3 \xi+4 M+1)}{\xi+M}, \tag{1.6}
\end{equation*}
$$

where $M=\sup _{\bar{\Omega}}|V| / \lambda_{1}$ and $\xi=\lambda_{2} / \lambda_{1}$.
Remark. If $V=0$ in (1.6), from (1.6), it is easy to see that

$$
\frac{\lambda_{2}+\lambda_{3}+\cdots+\lambda_{n+1}}{\lambda_{1}} \leq n+3+\frac{\lambda_{1}}{\lambda_{2}}
$$

(1.4) follows. Hence, (1.6) extends the inequality (1.4).

## 2. Proof of Theorem

Let $u_{i}$ be the orthonormal eigenvalue function with respect to $L^{2}$ inner product corresponding to $\lambda_{i}$, that is,

$$
\int_{\Omega} u_{i} u_{j}=\delta_{i j}, \quad \text { for any } i, j
$$

We choose rectangular coordinates $\widetilde{x}^{1}, \widetilde{x}^{2}, \ldots, \widetilde{x}^{n}$ of the Euclidean space $\mathbb{R}^{n}$ by taking as origin the center of gravity of $\Omega$ with mass-distribution $u_{1}^{2}$ such that

$$
\begin{equation*}
\int_{\Omega} \tilde{x}^{i} u_{1}^{2}=0, \quad \text { for } i=1,2, \cdots, n \tag{2.1}
\end{equation*}
$$

Defining an $n \times n$-matrix $B$ as follows:

$$
B:=\left(b_{i j}\right)
$$

where $b_{i j}=\int_{\Omega} \widetilde{x}^{i} u_{1} u_{j+1}$. Using the orthogonalization of Gram and Schmidt, we know that there exist an upper triangle matrix $R=\left(R_{i j}\right)$ and an orthogonal matrix $Q=\left(q_{i j}\right)$ such that $R=Q B$, that is,

$$
\begin{equation*}
R_{i j}=\sum_{k=1}^{n} q_{i k} b_{k j}=\int_{k} \sum_{k=1}^{n} q_{i k} \widetilde{x}^{k} u_{1} u_{j}=0, \quad 2 \leq j \leq i \leq n . \tag{2.2}
\end{equation*}
$$

Setting $x^{i}=\sum_{j=1}^{n} q_{i j} \tilde{x}^{j}$. From (2.1) and (2.2), we arrive at

$$
\begin{equation*}
\int_{\Omega} x_{i} u_{1} u_{j}=0, \quad \text { for } 1 \leq j \leq i \leq n \tag{2.3}
\end{equation*}
$$

Let $\varphi_{i}=x_{i} u_{1}$. Then $\varphi_{i}=0$ on $\partial \Omega$ and

$$
\int_{\Omega} \varphi_{i} u_{j}=0, \quad \text { for } 1 \leq j \leq i \leq n
$$

One gets from Rayleigh-Ritz inequality that

$$
\begin{equation*}
\lambda_{i+1} \leq \frac{\int_{\Omega} \varphi_{i}(-\Delta+V) \varphi_{i}}{\int_{\Omega} \varphi_{i}^{2}} \tag{2.4}
\end{equation*}
$$

Note that

$$
(-\Delta+V) \varphi_{i}=\lambda_{1} x_{i} u_{1}-2 u_{1, x_{i}}
$$

where $u_{1, x_{i}}=\partial u_{1} / \partial x_{i}$. It follows that

$$
\begin{align*}
\int_{\Omega} \varphi_{i}(-\Delta+V) \varphi_{i} & =\int_{\Omega} \varphi_{i}\left(\lambda_{1} x_{i} u_{1}-2 u_{1, x_{i}}\right) \\
& =\lambda_{1} \int_{\Omega} \varphi_{i}^{2}-2 \int_{\Omega} x_{i} u_{1} u_{1, x_{i}} \\
& =\lambda_{1} \int_{\Omega} \varphi_{i}^{2}-\int_{\Omega} x_{i}\left(u_{1}^{2}\right)_{, x_{i}}  \tag{2.5}\\
& =\lambda_{1} \int_{\Omega} \varphi_{i}^{2}+\int_{\Omega} u_{1}^{2} \\
& =\lambda_{1} \int_{\Omega} \varphi_{i}^{2}+1 .
\end{align*}
$$

(2.5) combining with (2.4) yields

$$
\begin{equation*}
\lambda_{i+1} \leq \lambda_{1}+\left(\int_{\Omega}\left(x_{i} u_{1}\right)^{2}\right)^{-1} . \tag{2.6}
\end{equation*}
$$

By integration by parts, it holds that

$$
\int_{\Omega} u_{1}^{\alpha+1}=-\int_{\Omega} x_{i}\left(u_{1}^{\alpha+1}\right)_{, x_{i}}=-(\alpha+1) \int_{\Omega}\left(x_{i} u_{1}\right)\left(u_{1}^{\alpha-1} u_{1, x_{i}}\right) .
$$

For $\alpha>1 / 2$, it follows from the Cauchy-Schwarz inequality that

$$
\begin{aligned}
\left(\int_{\Omega} u_{1}^{\alpha+1}\right)^{2} & =(\alpha+1)^{2}\left(\int_{\Omega}\left(x_{i} u_{1}\right)\left(u_{1}^{\alpha-1} u_{1, x_{i}}\right)\right)^{2} \\
& \leq(\alpha+1)^{2} \int_{\Omega}\left(x_{i} u_{1}\right)^{2} \int_{\Omega}\left(u_{1}^{\alpha-1} u_{1, x_{i}}\right)^{2} \\
& =\frac{(\alpha+1)^{2}}{2 \alpha-1} \int_{\Omega}\left(x_{i} u_{1}\right)^{2} \int_{\Omega}\left(u_{1}^{2 \alpha-1}\right)_{, x_{i}} u_{1, x_{i}} \\
& =\frac{-(\alpha+1)^{2}}{2 \alpha-1} \int_{\Omega}\left(x_{i} u_{1}\right)^{2} \int_{\Omega} u_{1}^{2 \alpha-1} u_{1, x_{i} x_{i}}
\end{aligned}
$$

Thus

$$
\begin{equation*}
\left(\int_{\Omega}\left(x_{i} u_{1}\right)^{2}\right)^{-1} \leq \frac{-(\alpha+1)^{2}}{2 \alpha-1} \frac{\int_{\Omega} u_{1}^{2 \alpha-1} u_{1, x_{i} x_{i}}}{\left(\int_{\Omega} u_{1}^{\alpha+1}\right)^{2}} \tag{2.7}
\end{equation*}
$$

Applying (2.7) to (2.6), one gets

$$
\begin{align*}
\frac{\lambda_{2}+\lambda_{3}+\cdots+\lambda_{n+1}}{\lambda_{1}} & \leq n+\frac{(\alpha+1)^{2}}{2 \alpha-1} A(\alpha) \sup _{\bar{\Omega}}\left(1-\frac{V}{\lambda_{1}}\right)  \tag{2.8}\\
& \leq n+(M+1) \frac{(\alpha+1)^{2}}{2 \alpha-1} A(\alpha),
\end{align*}
$$

where $A(\alpha)=\int_{\Omega} u_{1}^{2 \alpha} /\left(\int_{\Omega} u_{1}^{\alpha+1}\right)^{2}$.
In the following, we will find an upper bound of

$$
\begin{equation*}
\frac{(\alpha+1)^{2}}{2 \alpha-1} A(\alpha) \tag{2.9}
\end{equation*}
$$

Define

$$
\phi=u_{1}^{\alpha}-u_{1} \int_{\Omega} u_{1}^{\alpha+1}, \text { for } \alpha>1
$$

Then we have

$$
\int_{\Omega} \phi u_{1}=0
$$

This means that

$$
\begin{equation*}
\lambda_{2} \leq \frac{\int_{\Omega} \phi(-\Delta+V) \phi}{\int_{\Omega} \phi^{2}} \tag{2.10}
\end{equation*}
$$

Note that

$$
\begin{aligned}
\alpha \int_{\Omega} u_{1}^{\alpha-1}\left|\nabla u_{1}\right|^{2} & =\int_{\Omega} u_{1}^{\alpha}(-\Delta) u_{1} \\
& =\int_{\Omega}\left(\lambda_{1}-V\right) u_{1}^{\alpha+1}, \\
(2 \alpha-1) \int_{\Omega} u_{1}^{2 \alpha-2}\left|\nabla u_{1}\right|^{2} & =\int_{\Omega} u_{1}^{2 \alpha-1}(-\Delta) u_{1} \\
& =\int_{\Omega}\left(\lambda_{1}-V\right) u_{1}^{2 \alpha},
\end{aligned}
$$

and

$$
(-\Delta+V) \phi=-\alpha(\alpha-1) u_{1}^{\alpha-2}\left|\nabla u_{1}\right|^{2}+\left(\alpha \lambda_{1}-\alpha V+V\right) u_{1}^{\alpha}-\lambda_{1} u_{1} \int_{\Omega} u_{1}^{\alpha+1} .
$$

Hence, we have

$$
\begin{align*}
\int_{\Omega} \phi(-\Delta+V) \phi & =\frac{\alpha^{2}}{2 \alpha-1} \lambda_{1} \int_{\Omega} u_{1}^{2 \alpha}-\frac{(\alpha-1)^{2}}{2 \alpha-1} \int_{\Omega} V u_{1}^{2 \alpha}-\lambda_{1}\left(\int_{\Omega} u_{1}^{\alpha+1}\right)^{2}  \tag{2.11}\\
& \leq\left(\frac{\alpha^{2}}{2 \alpha-1}+\frac{(\alpha-1)^{2}}{2 \alpha-1} M\right) \lambda_{1} \int_{\Omega} u_{1}^{2 \alpha}-\lambda_{1}\left(\int_{\Omega} u_{1}^{\alpha+1}\right)^{2}
\end{align*}
$$

From (2.10) and (2.11), we arrive at

$$
\begin{equation*}
\frac{\lambda_{2}}{\lambda_{1}} \leq \frac{\left(\frac{\alpha^{2}}{2 \alpha-1}+\frac{(\alpha-1)^{2}}{2 \alpha-1} M\right) A(\alpha)-1}{A(\alpha)-1} \tag{2.12}
\end{equation*}
$$

Again, by using the Cauchy-Schwarz inequality, one gets

$$
\left(\int_{\Omega} u_{1}^{\alpha+1}\right)^{2}=\left(\int_{\Omega} u_{1}^{\alpha} u_{1}\right)^{2} \leq \int_{\Omega} u_{1}^{2 \alpha} \int_{\Omega} u_{1}^{2}=\int_{\Omega} u_{1}^{2 \alpha} .
$$

This means that $A(\alpha)>1$ for $\alpha>1$. If $\alpha$ is restricted to the condition

$$
\xi-\left(\frac{\alpha^{2}}{2 \alpha-1}+\frac{(\alpha-1)^{2}}{2 \alpha-1} M\right)>0
$$

that is

$$
\begin{equation*}
1<\alpha<\frac{(M+\xi)+\sqrt{(M+\xi)(\xi-1)}}{M+1} \tag{2.13}
\end{equation*}
$$

Then (2.12) is equivalent to

$$
\begin{equation*}
A(\alpha) \leq \frac{\xi-1}{\xi-\left(\frac{\alpha^{2}}{2 \alpha-1}+\frac{(\alpha-1)^{2}}{2 \alpha-1} M\right)} \tag{2.14}
\end{equation*}
$$

where $\xi=\lambda_{2} / \lambda_{1}$. Inserting (2.14) into (2.8) yields

$$
\begin{equation*}
\frac{\lambda_{2}+\lambda_{3}+\cdots+\lambda_{n+1}}{\lambda_{1}} \leq n+(M+1)(\xi-1) f(\alpha), \tag{2.15}
\end{equation*}
$$

where

$$
f(\alpha)=\frac{(\alpha+1)^{2}}{(2 \alpha-1) \xi-\left[\alpha^{2}+(\alpha-1)^{2} M\right]}
$$

The minimum of $f(\alpha)$ as a function of $\alpha$ in the range (2.13) is

$$
\frac{(M+1)(3 \xi+4 M+1)}{(\xi+M)(\xi-1)}
$$

and this is attained at

$$
\alpha=\frac{2 \xi+2 M}{\xi+2 M+1} .
$$

Hence, (2.15) yields

$$
\begin{aligned}
\frac{\lambda_{2}+\lambda_{3}+\cdots+\lambda_{n+1}}{\lambda_{1}} & \leq n+(M+1)(\xi-1) f\left(\frac{2 \xi+2 M}{\xi+2 M+1}\right) \\
& =n+\frac{(M+1)(3 \xi+4 M+1)}{\xi+M}
\end{aligned}
$$

This concludes the proof of theorem.

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